

An Inverse Problem for Two-Frequency Photon Transport in a Slab

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Interstellar cloud



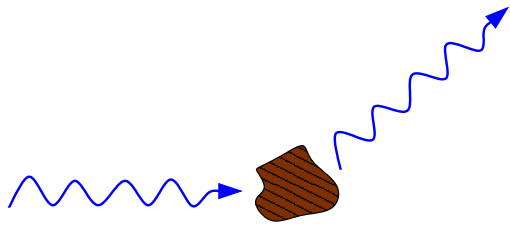
A large, gaseous, gravitational system (~ 10 ly) composed by atoms, simple and complex molecules, “dust” grains.

Scattering of light

Light coming from external or internal stars interacts with dust grains and undergoes three main processes:

Scattering of light

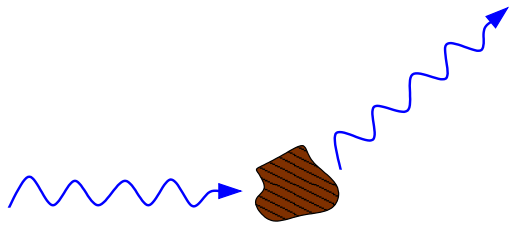
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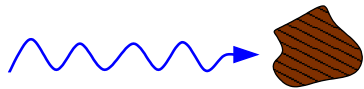
direction scattering

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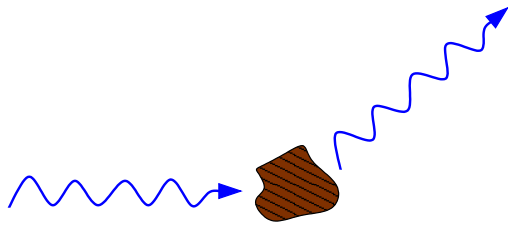
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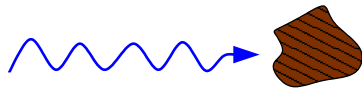
capture

Scattering of light

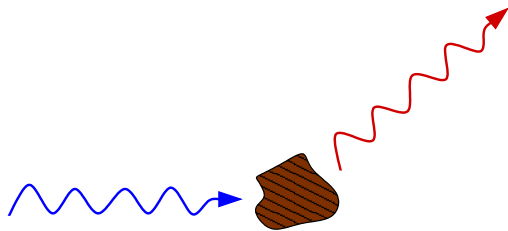
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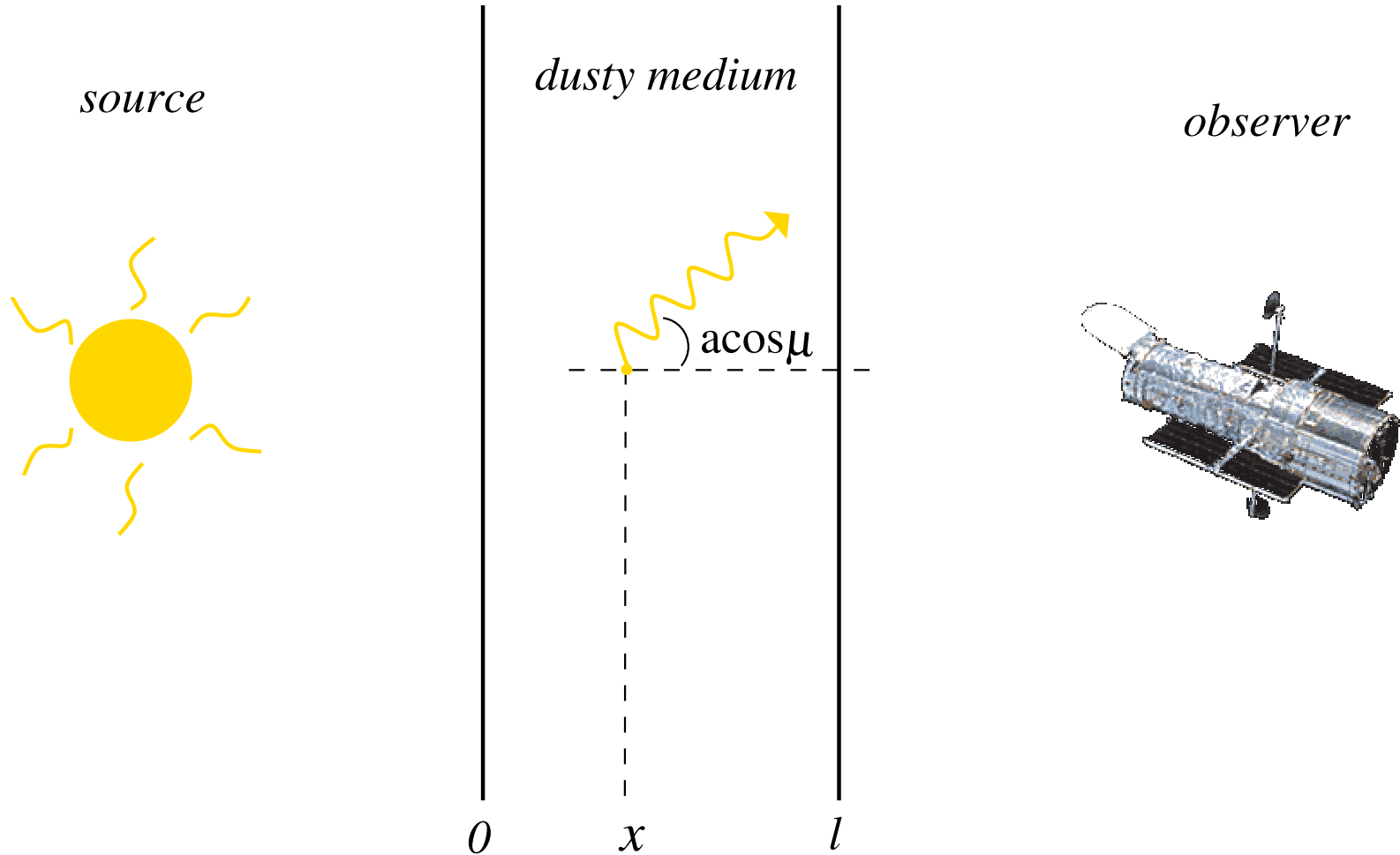


capture



direction/frequency scattering

Slab symmetry



Mathematical model

$f_1(x, \mu)$ = density at x of UV-photons with direction μ

$f_2(x, \mu)$ = density at x of IR-photons with direction μ

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$\Sigma_{i \rightarrow j}$ = macroscopic scattering cross-sections;

$p_{i \rightarrow i}(\mu' \rightarrow \mu)$ = transition probabilities;

$\Sigma_i := \Sigma_{i,c} + \Sigma_{i \rightarrow 1} + \Sigma_{i \rightarrow 2}$ = total cross-sections;

$\Sigma_{i,c}$ = capture cross-sections.

Mathematical model

Stationary radiative transfer equations:

$$\mu \frac{\partial f_1}{\partial x}(x, \mu) + \Sigma_1 f_1(x, \mu) = \Sigma_{1 \rightarrow 1} \int_{-1}^1 p_{1 \rightarrow 1}(\mu' \rightarrow \mu) f_1(x, \mu') d\mu'$$

$$\begin{aligned} \mu \frac{\partial f_2}{\partial x}(x, \mu) + \Sigma_2 f_2(x, \mu) &= \Sigma_{2 \rightarrow 2} \int_{-1}^1 p_{2 \rightarrow 2}(\mu' \rightarrow \mu) f_2(x, \mu') d\mu' \\ &+ \Sigma_{1 \rightarrow 2} \int_{-1}^1 p_{1 \rightarrow 2}(\mu' \rightarrow \mu) f_1(x, \mu') d\mu \end{aligned}$$

Dust densities

We assume that two kinds of dust, with unknown constant, densities ρ_1 and ρ_2 are present in the medium.

Introducing a *dust index* $d = 1, 2$ and denoting by $\sigma_{i \rightarrow j}^d$, σ_i^d , $\sigma_{i,c}^d$ the microscopic cross-sections for each kind of dust, we can write:

$$\Sigma_{i \rightarrow j} = \rho_1 \sigma_{i \rightarrow j}^1 + \rho_2 \sigma_{i \rightarrow j}^2$$

$$\Sigma_i = \rho_1 \sigma_i^1 + \rho_2 \sigma_i^2$$

$$\Sigma_{i,c} = \rho_1 \sigma_{i,c}^1 + \rho_2 \sigma_{i,c}^2$$

Inflow conditions

We assume that the *inflows* from the left

$$f_1(0, \mu) = \varphi_1^+(\mu), \quad f_2(0, \mu) = \varphi_2^+(\mu), \quad \text{for } \mu \in (0, 1),$$

and from the right

$$f_1(l, \mu) = \varphi_1^-(-\mu), \quad f_2(l, \mu) = \varphi_2^-(-\mu), \quad \text{for } \mu \in (-1, 0),$$

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The typical **direct** problem is finding the left and right *outflows*.

The inverse problem

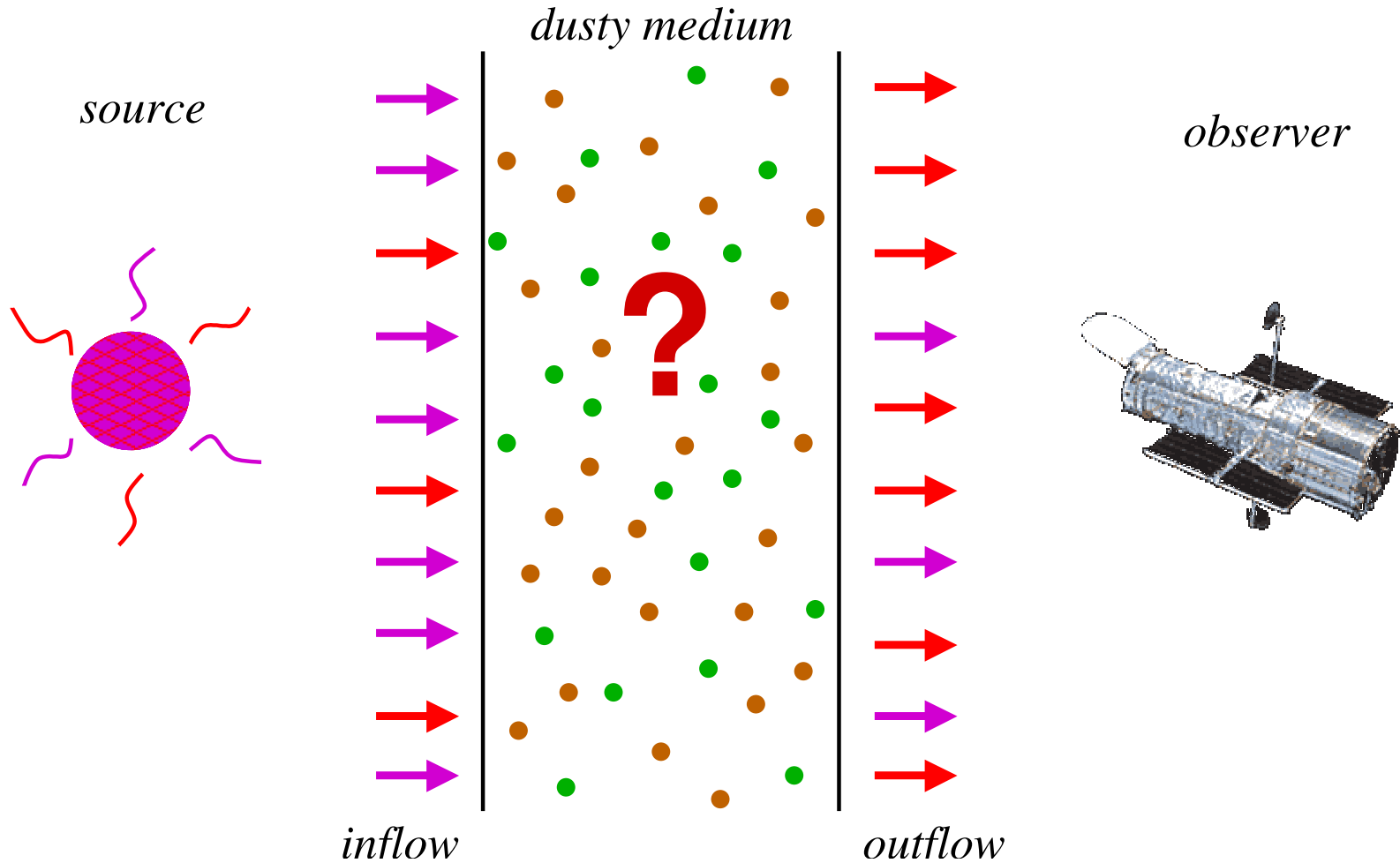
Problem: find the unknown dust densities ρ_1 and ρ_2 , assuming that the *integrated right outflows* at the two frequencies

$$H_1 := \int_0^1 f_1(l, \mu) \mu d\mu,$$

$$H_2 := \int_0^1 f_2(l, \mu) \mu d\mu,$$

are known.

The inverse problem



The main result

We make the following assumptions:

A1. *the frequency-scattering vanishes, i.e. $\Sigma_{1 \rightarrow 2} \equiv 0$;*

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A1. *the frequency-scattering vanishes, i.e. $\Sigma_{1 \rightarrow 2} \equiv 0$;*

A2. *the left-inflow data φ_1^+ and φ_2^+ are positive on nonzero-measure sets;*

A3. *$\sigma_{i,c}^j > 0$ for $i, j \in \{1, 2\}$ and*

$$\begin{vmatrix} \sigma_{1,c}^1 & \sigma_{1,c}^2 \\ \sigma_{2,c}^1 & \sigma_{2,c}^2 \end{vmatrix} \neq 0$$

The main result

THEOREM *Under assumptions **A1**, **A2**, **A3**, the mapping densities-to-outflows,*

$$(\rho_1, \rho_2) \mapsto (H_1, H_2),$$

is globally invertible and, therefore, the inverse problem is well-posed.

Sketch of the proof

By separating leftward and rightward photons

$$f_i^+(x, \mu) := f_i(x, \mu), \quad f_i^-(x, \mu) := f_i(l - x, -\mu)$$

(for $\mu \in (0, 1)$ and $i = 1, 2$), the direct problem can be recast into an “evolution” equation (from the inflow to the outflow):

$$\begin{cases} \frac{d}{dx} \mathbf{f}(x) = A \mathbf{f}(x), & x \in [0, l], \\ \mathbf{f}(0) = \varphi. \end{cases}$$

Sketch of the proof

where $\mathbf{f} := (f_1^+, f_1^-, f_2^+, f_2^-)$, $\boldsymbol{\varphi} := (\varphi_1^+, \varphi_1^-, \varphi_2^+, \varphi_2^-)$
and

$$A = \frac{1}{\mu} \begin{pmatrix} -\Sigma_1 + K_1^{++} & K_1^{-+} & 0 & 0 \\ K_1^{+-} & -\Sigma_1 + K_1^{--} & 0 & 0 \\ 0 & 0 & -\Sigma_2 + K_2^{++} & K_2^{-+} \\ 0 & 0 & K_2^{+-} & -\Sigma_2 + K_2^{--} \end{pmatrix}$$

($K_i^{s_1 s_2}$ are suitable scattering operators).

Sketch of the proof

By using the Lumer-Phillips generation problem we can prove that A generates a contraction semigroup (on the Banach space $L^1([0, 1], \mu d\mu)^4$), so that the inflow-to-outflow mapping is explicitly given by:

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Now, A is linear with respect to the dust densities:

$$A = \rho_1 A_1 + \rho_2 A_2$$

and assumption **A1** implies commutativity:

$$A_1 A_2 = A_2 A_1.$$

Sketch of the proof

From commutativity we obtain $e^{lA} = e^{l\rho_1 A_1} e^{l\rho_2 A_2}$, which gives

$$\frac{\partial f(l)}{\partial \rho_j} = l A_j e^{lA} \varphi.$$

This allows to prove that the integrated outflows

$$H_i = \int_0^1 f_i^+(l, \mu) \mu d\mu = \int_0^1 \left[e^{lA} \varphi \right]_i^+ (\mu) \mu d\mu$$

are continuously-differentiable with respect to ρ_j .

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$$\frac{\partial H_i}{\partial \rho_j} < 0, \quad i, j \in \{1, 2\}.$$

These conditions, together, imply the global invertibility of $(\rho_1, \rho_2) \mapsto (H_1, H_2)$. □

An inversion algorithm

A simple bisection-like algorithm can be used to solve the inverse problem:

1. by solving the direct problem, the range of the density-to-outflow mapping is divided into cells;

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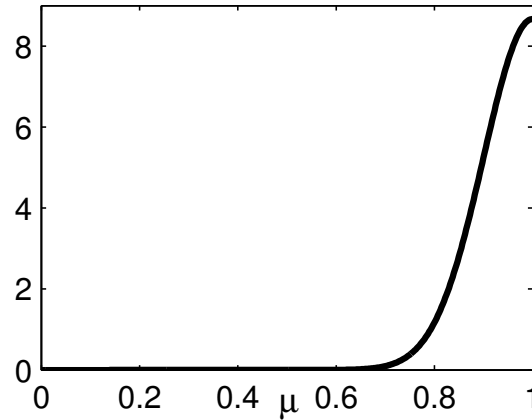
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2. the cell containing the measured value (H_1^0, H_2^0) is detected;
3. a refined grid is produced within that cell.
4. steps 2 and 3 are repeated up to desired order of accuracy.

Numerical experiments

Inflow datum:



Dust data:

λ (μm)	Dust kind	radius (μm)	σ_c (μm^2)	σ_s (μm^2)
0.1	graphite	0.25	0.18	0.29
10.0	graphite	0.25	0.93 E-2	0.25 E-3
0.1	silicate	1.00	2.67	4.03
10.0	silicate	1.00	4.22	0.50

Numerical experiments

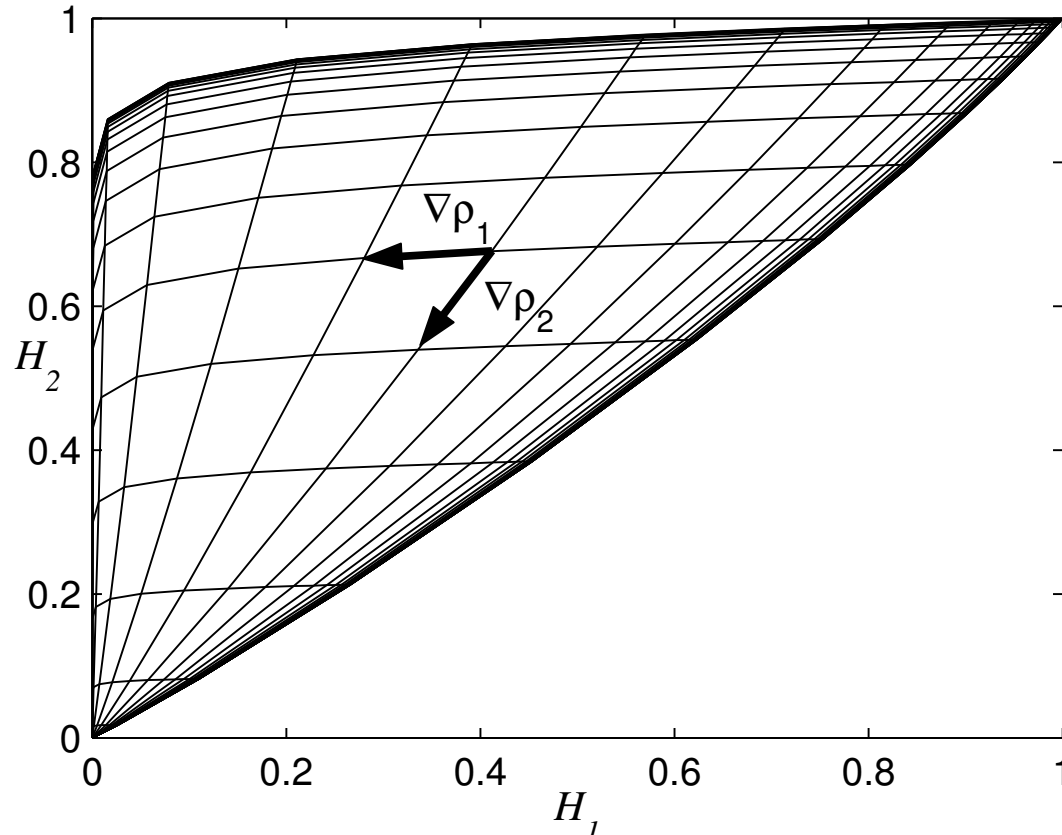
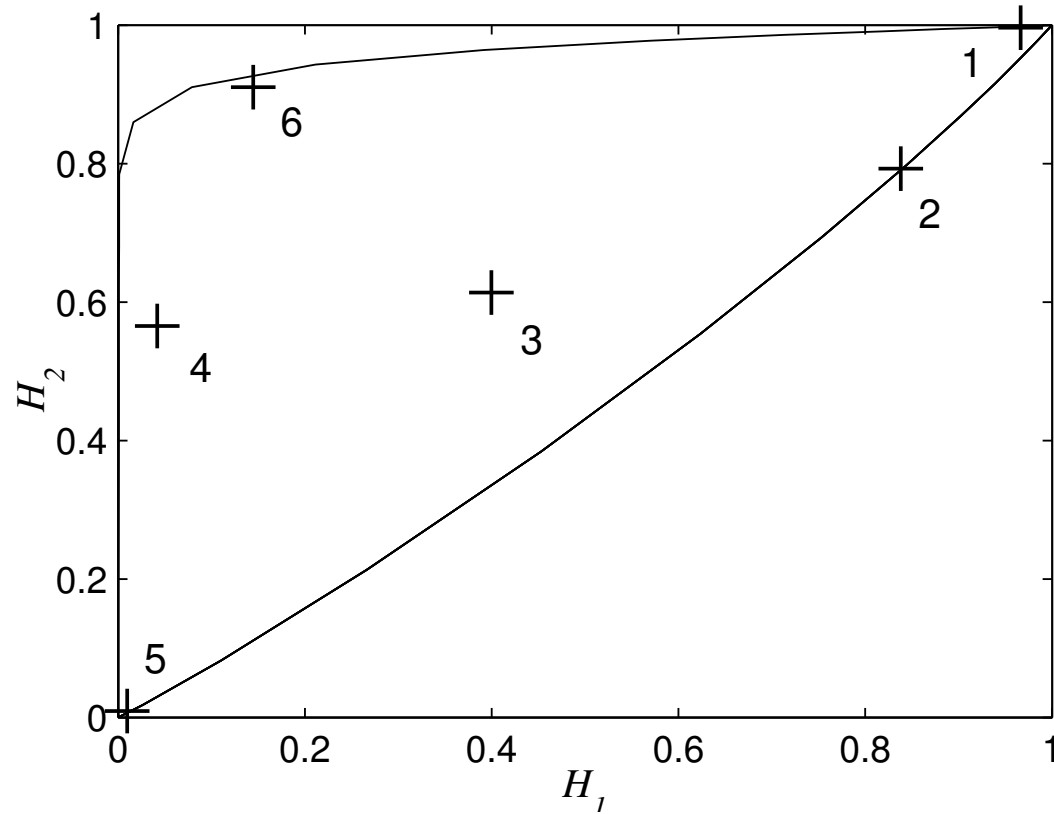


Image of the square $10^{-7} \text{ m}^{-3} \leq \rho_1, \rho_2 \leq 10^{-1} \text{ m}^{-3}$ under the mapping $(\rho_1, \rho_2) \mapsto (H_1, H_2)$. The grid is logarithmically spaced.

Numerical experiments

Test points:



Numerical experiments

	H_1^0	H_2^0	ρ_1^0	ρ_2^0	Δ_1	Δ_2	T
1	9.664 E-1	9.962 E-1	1.50 E-4	5.00 E-7	1.5 E-4	1.3 E-4	9.4 E+0
2	8.379 E-1	7.927 E-1	1.00 E-5	5.00 E-5	4.0 E-4	1.9 E-4	2.8 E+1
3	3.996 E-1	6.136 E-1	2.00 E-3	1.00 E-4	4.3 E-4	2.8 E-4	4.4 E+1
4	4.195 E-2	5.652 E-1	1.00 E-2	1.00 E-4	2.0 E-5	7.6 E-4	4.3 E+2
5	9.857 E-3	9.197 E-3	1.25 E-3	1.00 E-3	1.6 E-4	1.9 E-16	4.8 E+2
6	1.446 E-1	9.102 E-1	7.00 E-3	5.00 E-6	1.0 E-4	9.1 E-5	4.8 E+2

(H_1^0, H_2^0) = measured outflows;

(ρ_1^0, ρ_2^0) = computed dust densities;

(Δ_1, Δ_2) = relative difference with the true values of (ρ_1, ρ_2) ;

T = CPU time.

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- the problem is solvable if the two kinds of dust have different absorption properties on the two measured frequencies, and in absence of frequency-scattering;
- the latter assumption seems to be just technical and we hope to get rid of it in future works;
- a simple bisection algorithm gives good results for realistic values of the physical constants.