

ABOUT THE NOTION OF NON- T -RESONANCE AND APPLICATIONS TO TOPOLOGICAL MULTIPLICITY RESULTS FOR ODES ON DIFFERENTIABLE MANIFOLDS

LUCA BISCONTI AND MARCO SPADINI

ABSTRACT. By using topological methods, mainly the degree of a tangent vector field, we establish multiplicity results for T -periodic solutions of parametrized T -periodic perturbations of autonomous ODEs on a differentiable manifold M . In order to provide insights into the key notion of T -resonance, we consider the elementary situations $M = \mathbb{R}$ and $M = \mathbb{R}^2$. So doing, we provide more comprehensive analysis of those cases and find improved conditions.

2000 Mathematics Subject Classification: 34A09; 34C25; 34C40.

Key words: Periodic solutions, ordinary differential equations on manifolds, multiplicity.

1. INTRODUCTION

The search of multiplicity results for periodic solutions is a growing area of research. In this paper, we consider parametrized periodic perturbations of autonomous ODEs on differentiable manifolds and concentrate upon those results which are attainable through an intrinsically topological route. In doing this, we follow mainly [7] and [5] basing our argument on the interplay between global and local results (Theorem 3.5). A key notion for the “local” part of this approach is that of *ejecting set* or *point* (see Section 3 for a definition). Although it is sometimes possible to prove directly that the property of being ejecting holds for some points, usually the most practical way is through a non- T -resonance condition (see, e.g., [3, Ch. 7], a similar idea can be traced back to Poincaré, see [14] for an exposition).

We point out that many other approaches to multiplicity results for ordinary differential equations (especially of second order) have been pursued successfully. An exhaustive list of those is impossible to give here; we only mention [2, 13, 16, 17, 20] and references therein which we consider representative of different aspects and techniques.

The main purpose of this paper is to explore the concept of (non-) T -resonant zero of a vector field, providing an intuitive geometric grasp for it. We attempt to achieve this goal by considering the simplest possible situations, namely $M = \mathbb{R}$ and $M = \mathbb{R}^2$, for very smooth vector fields so that the implicit function theorem helps us make sense of the local shape of the set of T -periodic solutions. Indeed, in our investigation we come across a rather sharp description of the local structure of this set for scalar differential equations providing multiplicity/non-existence results (Proposition 4.5). In the two-dimensional case, we find an improved condition for an isolated zero of a vector field on \mathbb{R}^2 to be ejecting (Theorem 4.12).

Throughout this paper, we deal with the following parametrized differential equation

$$(1) \quad \dot{x} = g(x) + \lambda f(t, x), \quad \lambda \in [0, \infty),$$

where, throughout this section, $g: M \rightarrow \mathbb{R}^k$ and $f: \mathbb{R} \times M \rightarrow \mathbb{R}^k$ are continuous tangent vector fields on a boundaryless manifold $M \subseteq \mathbb{R}^k$ and f is T -periodic in the first variable. Our above mentioned multiplicity result (Theorem 3.5) gives a lower bound on the number of T -periodic solutions of (1) for any fixed f when $\lambda \geq 0$ is sufficiently small. The bound is based only on the properties of g . Roughly speaking, we prove that if $M \subseteq \mathbb{R}^k$ is closed and g has $n - 1$ non- T -resonant zeros with local index not adding up to the degree of g , then, fixed the forcing field f , there exists $\lambda_* > 0$ such that (1) has at least n geometrically distinct T -periodic solutions. Recall that two T -periodic solutions x and y of (1) are said to be *geometrically distinct* if their images do not coincide; in fact, x and y are not geometrically distinct if and only if there exists $\tau \in (0, T)$ such that $x(t + \tau) = y(t)$.

Our approach is theoretically-oriented but the kind of equations treated is common to many physical and engineering models. For instance, under appropriate conditions, differential-algebraic equations, which are a very well-known modeling and simulation tool for constrained systems, can be reduced to the form (1) see, e.g. [11]. In fact, the technique

used in the proof of our main multiplicity result, Theorem 3.5, is shared by a number different contexts. For example differential-algebraic equations have been considered in [1, 19] and constrained mechanical systems in [7].

We conclude this introduction noting that the pictures present in the paper have been produced by the popular interactive plotting package Gnuplot [8] using data generated by simple software written by one of the authors (M. Spadini) for the purpose of analyzing the set of starting points for scalar and two-dimensional parametrized equations.¹

2. NOTIONS AND PRELIMINARY RESULTS

Theorem 2.3 below is the main basis of the results discussed here. It is, deeply grounded on topological methods, mainly fixed point index theory and the degree of a tangent vector field. However, here, the reader does not need to know the details of the former theory which are completely hidden in the proof. The same is not true for the latter, though. Good references for it are [9, 15]. Here, we confine ourselves to reminding that, given a tangent vector field $v: M \rightarrow \mathbb{R}^k$, here $M \subseteq \mathbb{R}^k$ denote a boundaryless differentiable manifold, and an open subset $W \subseteq M$ we say that v is admissible on W if $v^{-1}(0) \cap W$ is compact. In this case it is well-defined an integer $\deg(v, W)$, called the degree or characteristic of v in W , which roughly speaking counts algebraically the number of zeros of v that lie in W .²

The Excision property allows the introduction of the notion of index of an isolated zero. This concept will play a crucial role throughout this paper. Given an isolated zero p_0 of a vector field v we put $i(v, p_0) := \deg(v, U)$ where U is an isolating neighborhood of p_0 . Actually, if v is C^1 at p_0 , then $v'(p_0)$ is an endomorphism of $T_{p_0}M$ (see e.g. [15]) and, if $v'(p_0)$ is also invertible, then $i(v, p_0) = \text{sign det } v'(p_0)$. The additivity property of the degree implies that if v is admissible in W and $p_1, \dots, p_n \in W$ are isolated zeros of v , then

$$\deg(v, W) = \sum_{j=1}^n i(v, p_j) + \deg(v, W \setminus \{p_1, \dots, p_n\}).$$

Let us now recall some notions relative to the set of T -periodic solutions of (1).

Many considerations throughout the paper depend on the following version of Ascoli's theorem.

Theorem 2.1. *Let Y be a subset of \mathbb{R}^k and B a bounded equicontinuous subset of $C([a, b], Y)$. Then B is totally bounded in $C([a, b], Y)$. In particular, if Y is closed, then B is relatively compact.*

We say that $(\lambda, x) \in [0, \infty) \times C_T(M)$ is a T -pair if $x(\cdot)$ is a solution of (1) corresponding to λ . If $\lambda = 0$ and x is constant, then (λ, x) is said to be *trivial*. One may have nontrivial solutions even when $\lambda = 0$, as it happens, for instance, when $T = 2\pi$, $M = \mathbb{R}^2$ and $g(x_1, x_2) = (-x_2, x_1)$, $x = (x_1, x_2)$; clearly any 2π -periodic f can be taken for this example.

Denote by \mathcal{X} the subset of $[0, \infty) \times C_T(M)$ of all the T -pairs. Well-known properties of differential equations imply that \mathcal{X} is closed. Hence, as a closed subset of a locally complete space, it is locally complete as well. The following technical remark from [7] is noteworthy:

Remark 2.2. *The space \mathcal{X} of all the T -pairs is locally totally bounded. Thus, being locally complete, \mathcal{X} is locally compact. Moreover, if we assume that M is a complete manifold, then any bounded subset of \mathcal{X} is actually totally bounded. As a consequence, when the manifold $M \subseteq \mathbb{R}^k$ is a closed subset, closed and bounded sets of T -pairs are compact.*

¹The version of this software, adapted for scalar equations, is available through its author web page http://www.dma.unifi.it/~spadini/Software/01_Uutilities/StartingPoints-1D, whereas the one for two-dimensional equations is still under development.

²In the sense that when the zeros of v are all non-degenerate, then the set $v^{-1}(0) \cap W$ is finite and

$$\deg(v, W) = \sum_{p \in v^{-1}(0) \cap W} \text{sign det } v'(p).$$

It is in order to point out here that this notion of degree should not be confused with the (Brouwer) degree of maps between oriented manifolds, with which it morally coincides in the special case when $M = \mathbb{R}^k$. Indeed the degree of a tangent vector field satisfies all the classical properties of the Brouwer degree: *Solution, Excision, Additivity, Homotopy Invariance, Normalization* etc. (see also [12]).

As in [4] we stipulate some conventions for the sake of simplicity. Accordingly, we regard every space as its image in the following diagram of closed embeddings:

$$(2) \quad \begin{array}{ccc} [0, \infty) \times M & \longrightarrow & [0, \infty) \times C_T(M) \\ \uparrow & & \uparrow \\ M & \longrightarrow & C_T(M). \end{array}$$

In particular, M will be identified with its image in $C_T(M)$ under the embedding which associates to any $p \in M$ the map $\hat{p} \in C_T(M)$ constantly equal to p . And also, M will be regarded as the slice $\{0\} \times M \subseteq [0, \infty) \times M$. Similarly, $C_T(M)$ will be regarded as $\{0\} \times C_T(M)$.

According to these identifications, given an open subset Ω of $[0, \infty) \times C_T(M)$, by $\Omega \cap M$ we mean the open subset of M given by all $p \in M$ such that the pair $(0, \hat{p})$ belongs to Ω . If U is an open subset of $[0, \infty) \times M$, then $U \cap M$ represents the open set $\{p \in M : (0, p) \in U\}$. We point out that with the above conventions, $g^{-1}(0)$ can be viewed as the set of trivial T -pairs.

We will heavily use the following result about the structure of the set \mathcal{X} of T -pairs of (1) that was proved in [6].

Theorem 2.3. *Let $f: \mathbb{R} \times M \rightarrow \mathbb{R}^k$ and $g: M \rightarrow \mathbb{R}^k$ be continuous tangent vector fields defined on a (boundaryless) differentiable manifold $M \subseteq \mathbb{R}^k$, f being T -periodic in the first variable. Let Ω be an open subset of $[0, \infty) \times C_T(M)$, and assume that $\deg(g, \Omega \cap M)$ is well defined and nonzero. Then there exists a connected set Γ of nontrivial T -pairs of (1) in Ω whose closure in $[0, \infty) \times C_T(M)$ intersects $g^{-1}(0) \cap \Omega$ and is not contained in any compact subset of Ω . In particular, if M is closed in \mathbb{R}^k and $\Omega = [0, \infty) \times C_T(M)$, then Γ is unbounded.*

Definition 2.1. *Let f and g be smooth and let $\mathcal{S} \subseteq [0, \infty) \times M$ be the topological subspace consisting of all pairs (λ, p) such that the maximal solution of the following Cauchy problem:*

$$(3) \quad \begin{cases} \dot{x} = g(x) + \lambda f(t, x), \\ x(0) = p, \end{cases}$$

is T -periodic. The elements of \mathcal{S} are called starting points. If $(\lambda, p) \in \mathcal{S}$ the point p will be informally referred to as an initial point of a T -periodic solution.

An important fact that follows from the proof of Theorem 2.3 in [6] is the following:

Remark 2.4. *Let $h: \mathcal{X} \rightarrow \mathcal{S}$ be the map which assigns to any T -pair (λ, x) the starting point $(\lambda, x(0))$. Observe that h is continuous, onto and, since f and g are smooth, it is also one to one. Furthermore, by the continuous dependence on initial data, we get the continuity of $h^{-1}: \mathcal{S} \rightarrow \mathcal{X}$. Thus h maps $\mathcal{X} \cap \Omega$ homeomorphically onto $\mathcal{S} \cap U$. A starting point (λ, p_0) is called trivial when $\lambda = 0$ and $p_0 \in g^{-1}(0)$. Notice that trivial T -pairs correspond to the trivial starting points under the homeomorphism h .*

Suppose that f and g are such that Cauchy problem (3) admits unique maximal solution $x(\lambda, p, \cdot)$ for any $p \in M$ and $\lambda \in [0, \infty)$. Given $\lambda \in [0, \infty)$ and $t \in \mathbb{R}$, we set $P_t^\lambda(p) := x(\lambda, p, t)$ whenever this makes sense. Consider, in particular $t = T$. Well-known properties of differential equations imply that the domain \mathcal{D} of the map $(\lambda, p) \mapsto P_T^\lambda(p)$ is an open subset (possibly empty) of $[0, \infty) \times M$. In particular, for any $\lambda \geq 0$, the slice $\mathcal{D}_\lambda := \{p \in M : (\lambda, p) \in \mathcal{D}\}$ is the (open) domain of the map $p \mapsto P_T^\lambda(p)$. Clearly, the set \mathcal{S} defined above is contained in \mathcal{D} and we have

$$\mathcal{S} = \bigcup_{\lambda \in [0, \infty)} \text{Fix}(P_T^\lambda),$$

where $\text{Fix}(P_T^\lambda)$ denotes the set of fixed points of the map $P_T^\lambda(\cdot): \mathcal{D}_\lambda \rightarrow M$.

The importance of the property highlighted in the remark will become apparent in later sections. In fact, in order to illustrate the meaning of the non- T -resonance condition (see next section) we will consider explicit examples in \mathbb{R} and \mathbb{R}^2 and plot numerically (a part of) the image under h of the set Γ of Theorem 2.3. To the same end, in the case $M = \mathbb{R}^k$ we will also apply the implicit function theorem to the function $\mathcal{F}: \mathcal{D} \subseteq [0, \infty) \times \mathbb{R}^k \rightarrow \mathbb{R}^k$ defined by

$$\mathcal{F}(\lambda, p) := P_T^\lambda(p) - p$$

whose zeros are exactly the starting points of (1).

3. EJECTING SETS AND MULTIPLICITY

In this section we gather some of the notions which are central to our approach to multiplicity results. The main references for this section are [7, 5]. We start with some purely topological facts and definitions.

Let Y be a metric space and X a subset of $[0, \infty) \times Y$. Given $\lambda \geq 0$, let X_λ be the slice $\{y \in Y : (\lambda, y) \in X\}$. Moreover, given a topological space S and two subsets A and B , with $A \subseteq B$, \overline{A}^B and \overline{A} will denote the closure of A in B and in S , respectively. Analogously, by $\text{Fr}_B(A)$ and by $\text{Fr}(A)$ we refer to the boundary of A relative to B and to S respectively. Finally, $\#Z$ denotes the cardinality of a set Z .

Lemma 3.1 ([5]). *Let Y be a metric space and let X be a locally compact subset of $[0, \infty) \times Y$. Assume K is a compact relatively open subset of the slice X_0 . Then, for any sufficiently small open neighborhood U of K in Y , there exists a positive number δ such that*

$$X \cap ([0, \delta] \times \text{Fr}(U)) = \emptyset.$$

Definition 3.1. *Let X be a subset of $[0, \infty) \times Y$. We say that $A \subseteq X_0$ is an ejecting set (for X) if it is relatively open in X_0 and there exists a connected subset of X which meets A and is not contained in X_0 .*

If a singleton $\{p_0\} \subseteq X_0$ is an ejecting set, the point p_0 will be called, with slight abuse of notation, an *ejecting point*. An important class of ejecting points will be presented below where the notion of non- T -resonant point of (1) is introduced. In fact, in the context of multiplicity results for T -periodic solutions, the sets X and Y above will be defined as $X = \mathcal{X}$ and $Y = C_T(M)$.

We will say that a point $p_0 \in g^{-1}(0)$ is T -resonant for (1) if g is C^1 in a neighborhood of p_0 and if the linearized problem (on $T_{p_0}M$)

$$\begin{cases} \dot{x} = g'(p_0)x, \\ x(0) = x(T), \end{cases}$$

which corresponds to $\lambda = 0$, admits nontrivial solutions.

Notice that a point $p_0 \in g^{-1}(0)$ is not T -resonant if $g'(p_0)$ (which maps $T_{p_0}M$ into itself, see e.g. [15]) has no eigenvalues of the form $\frac{2n\pi i}{T}$ with $n \in \mathbb{Z}$. Thus, in particular, p_0 is an isolated zero of g and $i(g, p_0) \neq 0$.

The following lemma from [7] will play an important role. For reasons of future reference we reproduce its proof with some small simplifications.

Lemma 3.2. *Let $g: M \rightarrow \mathbb{R}^k$ be a C^1 tangent vector field. If $p_0 \in g^{-1}(0)$ is not T -resonant, then for any sufficiently small neighborhood V of the constant function $\hat{p}_0 \equiv p_0$ in $C_T(M)$ there exists a real number $\delta_V > 0$ such that $[0, \delta_V] \times \text{Fr}(V)$ does not contain any T -pair of (1).*

Proof. Since the set \mathcal{X} of the T -pairs of (1) is locally compact (see Remark 2.2), there exists an open neighborhood W of \hat{p}_0 in $C_T(M)$ and $\mu > 0$ such that $([0, \mu] \times \overline{W}) \cap \mathcal{X}$ is compact.

We claim that $(\{0\} \times \overline{W}) \cap \mathcal{X} = \{(0, \hat{p}_0)\}$ when W is sufficiently small. To prove this claim, assume by contradiction that there exists a sequence of nontrivial T -pairs of the form $\{(0, x_n)\}_{n \in \mathbb{N}}$ converging to $(0, \hat{p}_0)$. Thus $|x_n(t) - p_0| \rightarrow 0$ uniformly in t as $n \rightarrow \infty$. In particular, $\lim_{n \rightarrow \infty} x_n(0) = p_0$. Put

$$p_n := x_n(0) \quad \text{and} \quad u_n := \frac{p_n - p_0}{|p_n - p_0|}.$$

We can assume that $u_n \rightarrow u \in T_{p_0}M$. Let $P_t: M \rightarrow M$ be the Poincaré t -translation operator associated with the equation $\dot{x} = g(x)$. Since g is C^1 , the map $P_t(\cdot)$ is differentiable. Define $\Phi: M \rightarrow \mathbb{R}^k$ as $\Phi(\xi) = \xi - P_T(\xi)$. Then, Φ is differentiable as well and $\Phi(p_n) = \Phi(p_0) = 0$. Thus,

$$\Phi'(p_0)u = \lim_{n \rightarrow \infty} \frac{\Phi(p_n) - \Phi(p_0)}{|p_n - p_0|} = 0.$$

On the other hand $\Phi'(p_0): T_{p_0}M \rightarrow \mathbb{R}^k$ is given by $\Phi'(p_0)v = v - P'_T(p_0)v$ for all $v \in T_{p_0}M$. As it is well known, the map $\alpha: t \mapsto P'_t(p_0)u$ satisfies the Cauchy problem

$$(4) \quad \begin{cases} \dot{\alpha}(t) = g'(p_0)\alpha(t) \\ \alpha(0) = u. \end{cases}$$

Since p_0 is not T -resonant, $\Phi'(p_0)u = \alpha(0) - \alpha(T) \neq 0$. This contradiction proves the claim.

Let us now complete the proof. Take W satisfying the above properties. We have that $\{(0, \hat{p}_0)\} = \overline{W} \cap \mathcal{X}_0$ is compact and it is, being isolated, relatively open in \mathcal{X}_0 . The assertion now follows from Lemma 3.1. \square

Lemma 3.2 might look a little trivial at first. In fact, if $M = \mathbb{R}^k$ and f is sufficiently regular, its assertion can be deduced by the implicit function theorem applied to the equation $\mathcal{F}(\lambda, p) = 0$, where \mathcal{F} is as in Section 2. However, Lemma 3.2 holds true on any differentiable manifold and even when f is only continuous.

Remark 3.3. Let p_0 be a non- T -resonant zero of g and let V be a neighborhood of p_0 in $C_T(M)$ as in Lemma 3.2. Put $\Omega = [0, \infty) \times V$. It is not difficult to show that $\deg(g, \Omega \cap M) = i(g, p_0) \neq 0$. Theorem 2.3 shows that there exists a connected set $\Gamma \subseteq \Omega$ of nontrivial T -pairs of (1) in whose closure in $[0, \infty) \times C_T(M)$ intersects p_0 and is not contained in any compact subset of Ω . Thus, a combination of Theorem 2.3 and Lemma 3.2 shows that p_0 is an isolated point of \mathcal{X}_0 which is ejecting. It also imply that, if $\delta_V > 0$ is as in Lemma 3.2, then $\pi_1(\Gamma) \supseteq [0, \delta_V]$. Here π_1 denotes the projection of $[0, \infty) \times C_T(M)$ onto the first factor.

Our approach to multiplicity of T -periodic solutions is based on the following abstract multiplicity result from [5]:

Theorem 3.4 ([5]). Let Y be a metric space and let X be a locally compact subset of $[0, \infty) \times Y$. Assume that X_0 contains n pairwise disjoint ejecting subsets, $n - 1$ of which are compact. Then there exists $\lambda_* > 0$ such that $\#X_\lambda \geq n$ for any $\lambda \in [0, \lambda_*)$.

We are now in a position to state our main multiplicity result (compare [5, Th. 3.7] and [18, Th. 4.6]).

Theorem 3.5. Let $g: M \rightarrow \mathbb{R}^k$ and $f: \mathbb{R} \times M \rightarrow \mathbb{R}^k$ be tangent to the closed boundaryless submanifold M of \mathbb{R}^k . Assume that g is C^1 , $g^{-1}(0)$ is compact and f is T -periodic in t . Assume also that

(i) There are $n - 1$ non- T -resonant zeros of g , p_1, \dots, p_{n-1} such that

$$(5) \quad \sum_{i=1}^{n-1} i(g, p_i) \neq \deg(g, M);$$

(ii) The unperturbed equation

$$(6) \quad \dot{x} = g(x)$$

does not admit unbounded connected sets of T -periodic solutions in $C_T(M)$.

Then, for $\lambda > 0$ sufficiently small, the Equation (1) admits at least n geometrically distinct T -periodic solutions.

Example 3.6. Take $M = \mathbb{R}$ and let $g(x) = \frac{x+x^2}{1+x^2}$ and $f(t, x) = \sin t$ with $T = 2\pi$. We have two non- T -resonant zeros, $p_0 = 0$ and $p_1 = -1$. This implies the existence of two geometrically independent T -periodic solutions for small values of $\lambda \geq 0$. Figure 1 shows the set of starting points near p_0 and p_1 .

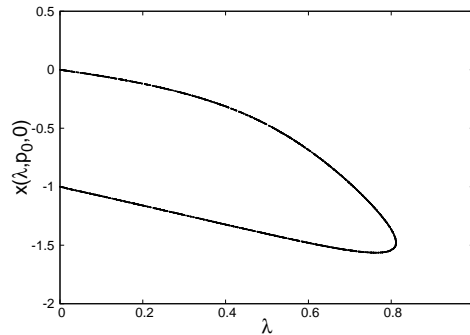


FIGURE 1. The set of starting points of Example 3.6.

Note that, when M is a compact boundaryless manifold condition (ii) of Theorem 3.5 holds automatically. Also, in this case, by the Poincaré-Hopf theorem one gets $\deg(g, M) = \chi(M)$, so that (5) becomes

$$\sum_{i=1}^n i(g, p_i) \neq \chi(M).$$

Thus we have the following consequence of Theorem 3.5:

Corollary 3.7. *Let $M \subseteq \mathbb{R}^k$ be a compact boundaryless manifold, and let f and g be as in Theorem 3.5. Assume that there are n non- T -resonant zeros of g , p_1, \dots, p_{n-1} such that*

$$\sum_{i=1}^{n-1} i(g, p_i) \neq \chi(M).$$

Then, for $\lambda > 0$ sufficiently small, the Equation (1) admits at least n geometrically distinct T -periodic solutions.

Example 3.8. *Take $M = S^2 \subseteq \mathbb{R}^3$ and $T = 2\pi$. Let $g: M \rightarrow \mathbb{R}^3$ be given by*

$$g(x, y, z) = \begin{cases} (-2yz^3, 2xz^3, 0) & \text{when } z \geq 0, \\ (xz^3, -yz^3, (y^2 - x^2)z^2) & \text{when } z \leq 0. \end{cases}$$

Clearly, g is C^1 and the points $p_1 := (0, 0, 1)$ and $p_2 = (0, 0, -1)$ are non- T -resonant zeros. with $i(g, p_1) = 1$ and $i(g, p_2) = -1$. So that $2 = \chi(S^2) \neq i(g, p_1) + i(g, p_2) = 0$. Therefore, for any 2π -periodic forcing term f Corollary 3.7 yields, 3 geometrically distinct 2π -periodic solutions of Equation (1) when $\lambda > 0$ is sufficiently small.

Proof of Theorem 3.5. Since p_1, \dots, p_{n-1} are non- T -resonant, we can find neighborhoods V_1, \dots, V_{n-1} such that

$$\overline{V_i} \cap g^{-1}(0) = \{p_i\} \quad \text{for } i = 1, \dots, n-1$$

By the definition of index of an isolated zero and the excision property, $\deg(g, V_i) = i(g, p_i)$, for $i = 1, \dots, n-1$. Define

$$V_0 = M \setminus \bigcup_{i=1}^{n-1} \overline{V_i}.$$

The additivity property of the degree yields

$$\deg(g, V_0) = \deg(g, M) - \sum_{i=1}^{n-1} i(p_i, g) \neq 0.$$

Define

$$\Omega = [0, \infty) \times C_T(V_0) \subset [0, \infty) \times C_T(M).$$

Theorem 2.3 implies that there exists an unbounded connected set Γ of nontrivial T -pairs of (1) in Ω whose closure in the space $[0, \infty) \times C_T(M)$ intersects $g^{-1}(0) \cap \Omega$. By assumption (ii) it follows that $g^{-1}(0) \cap V_0$ is an ejecting set of the set of T -pairs for (1). Remark 3.3 shows that p_1, \dots, p_{n-1} are ejecting sets (clearly compact). Thus, by Theorem 3.4, when $\lambda > 0$ is sufficiently small there are n different T -pairs, whence n different T -periodic solutions. That these are geometrically distinct follows from Lemma 3.2, for possibly smaller values of λ . \square

4. T -RESONANCE REVISITED

Going over the first part of the proof of Lemma 3.2 we see that what we actually prove is that if p is not an accumulation point of $\Gamma \cap (\{0\} \times M)$. This means that Γ does not intersect $\{0\} \times M$ infinitely many times in a neighborhood of $\{0\} \times \{p\}$, yet many pathological scenarios are possible. Something more can be said with the aid of the implicit function theorem if we assume some more regularity of f . In fact, we have to assume that $M = \mathbb{R}^n$ (this is harmless since we are going to make a local analysis) and allow λ to vary in \mathbb{R} (a whole neighborhood of 0) instead than in $[0, \infty)$.

Take $M = \mathbb{R}^n$ and assume that f and g are C^1 . As in Section 2, we set

$$\mathcal{F}(\lambda, p) = x(\lambda, p, T) - p$$

where $x(\lambda, p, \cdot)$ denotes the unique maximal solution of (3) and $\lambda \in \mathbb{R}$. Clearly, a pair $(\lambda, p) \in [0, \infty) \times M$ is a starting point (see Remark 2.4) if and only if $\mathcal{F}(\lambda, p) = 0$ and, also, starting points correspond to T -pairs when f and g are C^1 .

Let $p_0 \in M$ be a non- T -resonant zero of g . Then, simple computations show that, for all $v \in \mathbb{R}^n = T_{p_0}M$,

$$(7) \quad \partial_1 \mathcal{F}(0, p_0) = e^{Tg'(p_0)} \int_0^T e^{-sg'(p_0)} f(s, p_0) ds;$$

$$(8) \quad \partial_2 \mathcal{F}(0, p_0)v = \left(e^{Tg'(p_0)} - I \right) v.$$

Thus, since p_0 is non- T -resonant, $\partial_2 \mathcal{F}(0, p_0)$ is nonsingular and, by the implicit function theorem, there exists a function $p : \lambda \mapsto p(\lambda)$, defined in a neighborhood U of 0, such that $p(0) = 0$ and $\mathcal{F}(\lambda, p(\lambda)) \equiv 0$, so that all pairs $(\lambda, p(\lambda))$, $\lambda \in U$, are starting points. Clearly,

$$p'(0) = - \left(e^{Tg'(p_0)} - I \right)^{-1} \int_0^T e^{(T-s)g'(p_0)} f(s, p_0) ds$$

is well-defined. This means, geometrically, that when p_0 is non- T -resonant Γ crosses transversally $\{0\} \times M$ at p_0 . Conversely, however, we cannot conclude that this does not happen when p_0 is T -resonant (we shall see this in Example 4.4). To investigate this case we set $M = \mathbb{R}$. Later, we will consider the case $M = \mathbb{R}^2$ and, following the lead of Lemma 3.2, seek a second-order sufficient condition for an isolated zero with nonzero index to be ejecting.

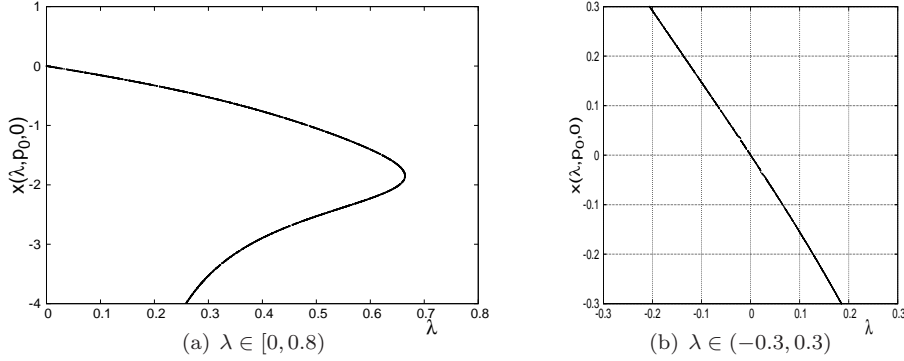


FIGURE 2. The set of starting points for Example 4.1.

Example 4.1. Let $M = \mathbb{R}$ and $g(x) = \frac{x}{1+x^2}$ and $f(t, x) = 1 + \cos(x + t)$ with $T = 2\pi$. Clearly g has exactly one zero $p_0 = 0$ which is non- T -resonant. Thus the implicit function theorem yields exists a function $p : \lambda \mapsto p(\lambda)$, defined in a neighborhood of 0, such that $p(0) = 0$ and the pairs $(\lambda, p(\lambda))$ are starting points for (1) for small values of $|\lambda|$. Figure 2(a) shows a portion of the set \mathcal{S} of starting points for the above choices of g and f . A simple computation yields $p'(0) = -3/2$ as it is illustrated by Figure 2(b).

4.1. The case $M = \mathbb{R}$. In general, Remark 3.3 shows that if $p_0 \in M$ is a non- T -resonant zero of g then the index of p_0 is nonzero and p_0 is ejecting. When $M = \mathbb{R}$ all isolated zeros of g with nonzero index are necessarily ejecting, as shown below.

Lemma 4.2. Let $M = \mathbb{R}$ and let p_0 be an isolated zero of g such that $i(g, p_0) \neq 0$. Then p_0 is an ejecting point for the set \mathcal{X} of T -pairs of (1). Consequently, p_0 is an ejecting point for the set \mathcal{S} of starting points of (1).

Sketch of the proof. As in the proof of Lemma 3.2, since the set \mathcal{X} of the T -pairs of (1) is locally compact, there exists an open neighborhood W of p_0 in $C_T(\mathbb{R})$ and $\mu > 0$ such that $\mathcal{X} \cap ([0, \mu] \times \overline{W})$ is compact.

Since T -periodic solutions of an autonomous equation on \mathbb{R} are necessarily constant and p_0 is an isolated zero, we see that $(\{0\} \times \overline{W}) \cap \mathcal{X} = \{(0, p_0)\}$. The proof is concluded in the same way as the proof of Lemma 3.2. \square

Let $p_0 \in \mathbb{R}$ be an isolated zero of g . Assume that p_0 is T -resonant, thus $e^{Tg'(p_0)} = 1$, and, since $M = \mathbb{R}$, $g'(p_0) = 0$. In particular $\partial_2 \mathcal{F}(0, p_0) = 0$ and $\partial_1 \mathcal{F}(0, p_0) = \int_0^T f(s, p_0) ds$. If we suppose that $\partial_1 \mathcal{F}(0, p_0) \neq 0$, the implicit function theorem yields a function $\lambda : p \mapsto \lambda(p)$, defined in a neighborhood V of p_0 , such that $\lambda(p_0) = 0$ and $\mathcal{F}(\lambda(p), p) \equiv 0$ for $p \in V$. Thus, in a neighborhood of p_0 , the set Γ of Theorem 2.3 is the graph of the map $\lambda : p \mapsto (\lambda(p), p)$, where $(\lambda(p), p)$ is regarded as a T -pair according to diagram (2). Geometrically, since

$$\lambda'(p_0) = -(\partial_1 \mathcal{F}(0, p_0))^{-1} \partial_2 \mathcal{F}(0, p_0) = -\frac{e^{Tg'(p_0)} - 1}{\int_0^T f(s, p_0) ds} = 0,$$

this means that Γ is tangent to $\{0\} \times \mathbb{R}$ at p_0 .

Example 4.3. Let $g(x) = \frac{x^3}{1+x^2}$ and $f(t, x) = 1 + \sin t$ with $T = 2\pi$. We have one T -resonant zero $p_0 = 0$, but $\int_0^T f(s, p_0) ds \neq 0$. Thus the set \mathcal{S} of starting points about 0 is tangent to the axis $\lambda = 0$, as shown in Figure 3.

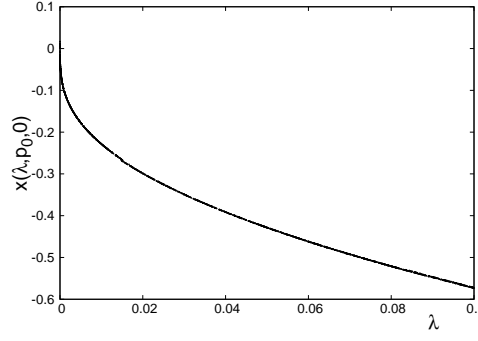


FIGURE 3. The set of starting points of Example 4.3.

Notice that nothing can be deduced through the implicit function theorem if we have both $g'(p_0) = 0$ and $\int_0^T f(s, p_0) ds = 0$. An example in this sense is provided below.

Example 4.4. Let $g(x) = \frac{x^3+x^2}{1+x^2}$ and $f(t, x) = \sin(t+x)$, $T = 2\pi$. We have one T -resonant $p_1 = 0$ and one non- T -resonant zero $p_2 = -1$, respectively. Since $\int_0^T f(s, p_1) ds = 0$. So the shape of the set \mathcal{S} of starting points about p_1 cannot be determined by the method above. Figure 4(a) shows a portion of the set of starting points for (1) when g and f are selected as above. Allowing λ in a complete neighborhood of 0 lets us understand the structure (see Figure 4(b)).

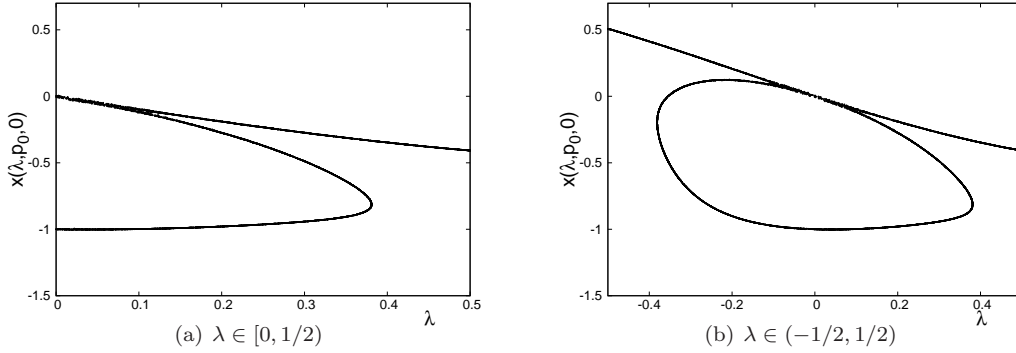


FIGURE 4. The set of starting points of Example 4.4.

We get more information about the shape of Γ by computing the second derivative of the function λ .

Let us stay with the assumption that p_0 is T -resonant, so that we have $g'(p_0) = 0$, and assume in addition that f and g are C^2 functions. As above, we suppose that $\int_0^T f(s, p_0) ds$ is not zero. Recall that for a C^2 function $\varphi: D \subseteq \mathbb{R} \rightarrow \mathbb{R}$, D open, we have

$$\varphi''(\xi) = \lim_{h \rightarrow 0} \frac{\varphi(\xi + hv) - 2\varphi(\xi) + \varphi(\xi - hv)}{h^2},$$

for any $\xi \in D$. Then we have

$$\begin{aligned} \partial_{22}x(0, p_0, t) &= \\ (9) \quad &= \lim_{h \rightarrow 0} \frac{1}{h^2} \int_0^t \left[g(x(0, p_0 + h, s)) - 2g(x(0, p_0, s)) + g(x(0, p_0 - h, s)) \right] ds \\ &= \int_0^t \left[\lim_{h \rightarrow 0^+} \frac{g'(p_0) \partial_{22}x(0, p_0, t) h^2 + g''(p_0) [\partial_2 x(0, p_0, s) h]^2 + o(h^2)}{h^2} \right] ds, \end{aligned}$$

and, since $g'(p_0) = 0$, we get

$$\partial_{22}x(0, p_0, t) = \int_0^t g''(p_0) [\partial_2 x(0, p_0, s)]^2 ds$$

Thus, setting $\alpha(t) = \partial_2 x(0, p_0, t)$ and $\beta = \partial_{22}x(0, p_0, t)$, we have that β solves the following Cauchy problem:

$$\begin{cases} \dot{\beta}(t) = g''(p_0) \alpha(t)^2, \\ \beta(0) = 0. \end{cases}$$

It is well-known that α solves (4) with $u = 1$. Since $g'(p_0) = 0$, we get $\alpha(t) \equiv 1$, so that $\partial_{22}\mathcal{F}(0, p_0) = \partial_{22}x(0, p_0, t) = Tg''(p_0)$. In order to determine $\lambda''(p_0)$ we differentiate twice the identity $\mathcal{F}(\lambda(p), p) = 0$ at p_0 :

$$\partial_{12}\mathcal{F}(0, p_0)\lambda'(p_0) + \partial_1\mathcal{F}(0, p_0)\lambda''(p_0) + \partial_{22}\mathcal{F}(0, p_0) = 0.$$

Recalling again that p_0 is T -resonant we have $\lambda'(p_0) = 0$. The above relation, along with (7), yields

$$(10) \quad \lambda''(p_0) = -\frac{g''(p_0)}{\frac{1}{T} \int_0^T f(s, p_0) ds}.$$

Suppose that $g''(p_0) \neq 0$. If the signs of $\int_0^T f(s, p_0) ds$ and $g''(p_0)$ disagree, then $\lambda''(p_0) > 0$. This implies that λ is a convex function in a neighborhood of p_0 . This fact has the following interesting consequence:

Proposition 4.5. *Let $M = \mathbb{R}$ and let p_0 be an isolated T -resonant zero of g . Assume that f and g are C^2 functions. Then,*

- (i) *Let $g''(p_0) \int_0^T f(s, p_0) ds > 0$. Then, given $\varepsilon > 0$, there exists $\lambda_* > 0$ such that Equation (1) has no T -periodic solutions that lie in $[p_0 - \varepsilon, p_0 + \varepsilon]$ for $\lambda \in [0, \lambda_*]$;*
- (ii) *Let $g''(p_0) \int_0^T f(s, p_0) ds < 0$. Assume that f has a separated variable form, that is there exist functions φ and h such that $f(t, x) = \varphi(t)h(x)$ for all $(t, x) \in \mathbb{R} \times \mathbb{R}$. Suppose also that T is the minimal period of φ . Then, given $\varepsilon > 0$, there exists $\lambda_* > 0$ such that Equation (1) has at least two geometrically distinct T -periodic solutions that lie in $[p_0 - \varepsilon, p_0 + \varepsilon]$ for $\lambda \in [0, \lambda_*]$.*
- (iii) *If $g''(p_0) \int_0^T f(s, p_0) ds \neq 0$ then $i(g, p_0) = 0$;*

Proof. For the sake of simplicity we shall assume throughout this proof that $p_0 = 0$.

(i) The argument above shows that there exists a function $\lambda : p \mapsto \lambda(p)$, defined in a neighborhood of $p_0 = 0$, such that $\lambda(0) = 0$ and $(\lambda(p), p)$ is a starting point that we regard as a T -pair according to diagram (2). As already observed, the T -resonance of 0 implies $g'(0) = 0$ so that we have $\lambda'(0) = 0$. Also, Formula (10) yields $\lambda''(0) < 0$. Since $\lambda(\cdot)$ is a C^2 function we can write

$$(11) \quad \lambda(p) = \frac{1}{2}\lambda''(0)p^2 + \sigma(|p|)p^2,$$

where σ is a continuous function with $\sigma(0) = 0$. Fix any $\bar{\lambda} > 0$. We have that, for a sufficiently small $\varepsilon > 0$, equation $\lambda(p) = \bar{\lambda}$ has no solution in the interval $[-\varepsilon, \varepsilon]$. In fact, if $p_{\bar{\lambda}} \in [-\varepsilon, \varepsilon]$ is such a solution, we have $p_{\bar{\lambda}} \neq 0$ and

$$(12) \quad \frac{1}{2}\lambda''(0) = \frac{\bar{\lambda}}{p_{\bar{\lambda}}^2} - \sigma(|p_{\bar{\lambda}}|).$$

Since $\sigma(0) = 0$ the right-hand-side of (12) is positive when $|p_{\bar{\lambda}}| < \varepsilon$ is sufficiently small, contradicting $\lambda''(0) < 0$. This proves the assertion.

(ii) Clearly, Equation (11) still holds but now we have $\lambda''(0) > 0$. Since $\sigma(0) = 0$, taking $\varepsilon > 0$ sufficiently small we can assume that

$$\alpha(p) := \frac{1}{2}\lambda''(0) + \sigma(|p|) > 0, \quad \forall p \in [-\varepsilon, \varepsilon].$$

Also, our assumption implies that $h(0) \neq 0$. Thus, restricting ε if necessary, we can assume that $h(p) \neq 0$ for all $p \in [-\varepsilon, \varepsilon]$.

We claim that the following equation:

$$\alpha(p)p^2 = \bar{\lambda}$$

admits two different solutions when $\bar{\lambda}$ is sufficiently small. To prove the claim put

$$\ell = \min \left\{ \max_{p \in [-\varepsilon, \varepsilon]} \{p\sqrt{\alpha(p)}\}, \max_{p \in [-\varepsilon, \varepsilon]} \{-p\sqrt{\alpha(p)}\} \right\},$$

which is clearly positive. The intermediate value theorem shows that the two equations

$$p\sqrt{\alpha(p)} = \lambda, \quad -p\sqrt{\alpha(p)} = \lambda$$

have both solutions (obviously different) in $[-\varepsilon, \varepsilon]$ for $\lambda \in [0, \sqrt{\ell}]$. Setting $\bar{\lambda} = \sqrt{\ell}$ we have the claim. Let us denote these two solutions by $p_1(\lambda)$ and $p_2(\lambda)$. Clearly, these points are initial points of T -periodic solutions of (1) which, for simplicity we denote by ξ_λ and η_λ . That is, for all $t \in \mathbb{R}$, we put

$$\xi_\lambda(t) = x(\lambda, p_1(\lambda), t) \quad \text{and} \quad \eta_\lambda(t) = x(\lambda, p_2(\lambda), t).$$

We need to prove that, there exists $\lambda_* \in (0, \bar{\lambda}]$ such that, when $\lambda \in (0, \lambda_*]$ the points ξ_λ and η_λ are geometrically distinct. By construction, both $p_1(\lambda)$ and $p_2(\lambda)$ tend to $p_0 = 0$ as $\lambda \rightarrow 0$. The continuous dependence on initial data implies that there exists $\lambda_* > 0$ such that the images of ξ_λ and η_λ are contained in $[-\varepsilon, \varepsilon]$. Suppose now, by contradiction, that ξ and η are not geometrically distinct. Then there exists $\tau \in (0, T)$ such that

$$\xi_\lambda(t) = x(\lambda, p_1(\lambda), t) = x(\lambda, p_2(\lambda), t + \tau) = \eta_\lambda(t + \tau).$$

Differentiating this identity with respect to t we get $\dot{\xi}_\lambda(t) = \dot{\eta}_\lambda(t + \tau)$ and, taking (1) into account, we get

$$\begin{aligned} g(\xi_\lambda(t)) + \lambda\varphi(t)h(\xi_\lambda(t)) &= \\ &= g(\eta_\lambda(t + \tau)) + \lambda\varphi(t + \tau)h(\eta_\lambda(t + \tau)) = g(\xi_\lambda(t)) + \lambda\varphi(t + \tau)h(\xi_\lambda(t)). \end{aligned}$$

This implies $\varphi(t)h(\xi_\lambda(t)) = \varphi(t + \tau)h(\xi_\lambda(t))$, for all t . Since h is never zero on $\xi_\lambda(t)$, we have $\varphi(t) = \varphi(t + \tau)$, $t \in \mathbb{R}$, which contradicts the minimal period assumption.

(iii) Suppose by contradiction that $i(g, p_0) \neq 0$. Assume first that $g''(p_0)$ and $\int_0^T f(s, p_0) ds$ are nonzero and their signs agree. On one hand, proceeding as in (i) we see that there are no T -periodic solutions in a neighborhood of p_0 for $\lambda > 0$ small. On the other hand we know that isolated zeros with nonzero index are ejecting for (1). This yields a contradiction. Assume now that $g''(p_0) \int_0^T f(s, p_0) ds < 0$. The same argument applied to equation

$$\dot{x} = g(x) - \lambda f(t, x)$$

(now f is replaced by $-f$) yields a contradiction again. \square

Example 4.6. Take $g(x) = \frac{x^3(1+x)^2(x-1)^2}{1+x^6}$ and $f(t, x) = \sin(t) + 1$, $T = 2\pi$. We have three T -resonant zeros, say $p_1 = 0$, $p_2 = 1$ and $p_3 = -1$, but $i(g, p_1) \neq 0$. So p_1 is ejecting. Also we have $\lambda''(p_1) = 0$, $\lambda''(p_2) = -4$, $\lambda''(p_3) = +4$. Thus there are 3 solutions for small $\lambda \geq 0$. Figure 5 shows a portion of the set of starting pairs for (1) when g and f are selected as above.

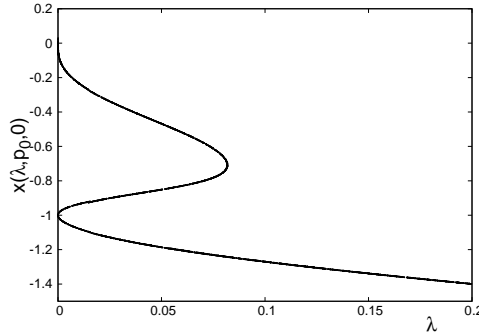


FIGURE 5. Starting points for Example 4.6.

Remark 4.7. Suppose p_0 is a T -resonant zero of g and $\int_0^T f(s, p_1) ds \neq 0$. The argument of Proposition 4.5 shows that, if the signs of $g''(p_0)$ and $\int_0^T f(s, p_1) ds$ disagree, then there are two T -periodic solutions near p_0 for small λ , but none if they agree. Figure 6 illustrates this phenomenon by showing a portion of the set of starting pairs for (1) when $g(x) = \frac{x^3+x^2}{1+x^2}$ and f is selected in different ways.

4.2. The case $M = \mathbb{R}^2$. When $M = \mathbb{R}^2$ it is no longer true that any isolated zero with nonzero index must be ejecting. We have, however, some structure:

Lemma 4.8. Let $M = \mathbb{R}^2$ and let p_0 be an isolated zero of g that is accumulation point of initial points of T -periodic orbits of $\dot{x} = g(x)$. Then, if r denotes a half-line with origin in p_0 , there exists a sequence $\{p_n\}_{n \in \mathbb{N}} \subseteq r$, with $p_n \rightarrow p_0$, such that the solution of $\dot{x} = g(x)$ starting from p_n at time $t = 0$ is T -periodic.

Sketch of the proof. Let $\{q_n\}_{n \in \mathbb{N}}$ be a sequence of initial points of T -periodic orbits of $\dot{x} = g(x)$ such that $q_n \rightarrow p_0$. Let x_n , $n \in \mathbb{N}$, be the maximal solution of the following Cauchy problems:

$$\begin{cases} \dot{x} = g(x), \\ x(0) = q_n. \end{cases}$$

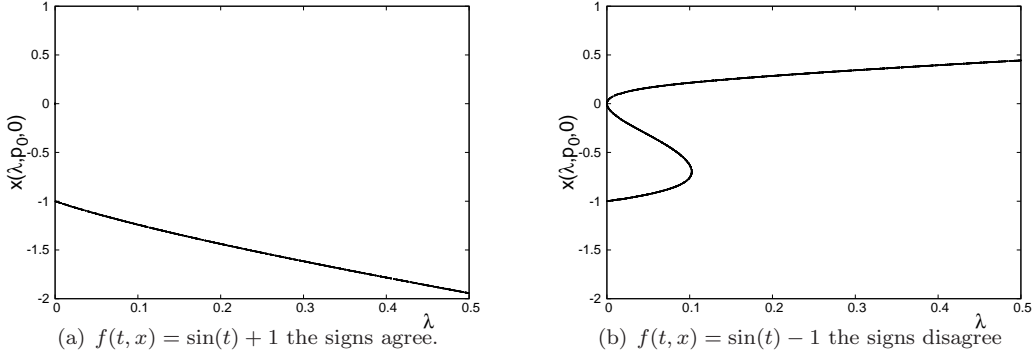


FIGURE 6. Two different cases for Example 4.7.

Then x_n converges uniformly on $[0, T]$ to the function $\hat{p}_0(t) \equiv p_0$. Let \mathcal{O}_n be the orbit that contains q_n , that is $\mathcal{O}_n = x_n([0, T])$. It is well-known that there exists z_n belonging to the part of plane enclosed by \mathcal{O}_n such that $g(z_n) = 0$ (see e.g. [10, Th. 2, §5, Ch. 11]). Clearly $z_n \rightarrow p_0$ but, since by assumption p_0 is an isolated zero of g , the only possibility is that $z_n = p_0$ eventually. This implies that the orbits \mathcal{O}_n enclose p_0 when n is sufficiently large. Thus, any half line with origin in p_0 necessarily meets \mathcal{O}_n . Let us choose a point p_n in $r \cap \mathcal{O}_n$. Clearly, since the equation is autonomous p_n is an initial point of a T -periodic solution (like all the points of \mathcal{O}_n) and $p_n \rightarrow p_0$. \square

We have the following necessary condition for an isolated zero to be accumulation point of initial points of T -periodic orbits of $\dot{x} = g(x)$.

Theorem 4.9. *Let $M = \mathbb{R}^2$ and let p_0 be an isolated zero of g that is accumulation point of initial points of T -periodic orbits of $\dot{x} = g(x)$. Then,*

$$(13) \quad \int_0^T e^{(T-s)g'(p_0)} g''(p_0) \left[e^{sg'(p_0)} v, e^{sg'(p_0)} v \right] ds = 0,$$

for all $v \in \mathbb{R}^2$.

The proof is based on the following technical result:

Lemma 4.10. *Let $M = \mathbb{R}^2$ and let p_0 be an isolated zero of g that is accumulation point of initial points of T -periodic orbits of $\dot{x} = g(x)$. Then,*

$$\partial_{22}x(0, p_0, T)[v, v] = 0,$$

for all $v \in \mathbb{R}^2$.

Proof. Let w be any unit vector in \mathbb{R}^2 and consider the half-line r_w generated by w with origin in $(0, 0)$, namely the set $r_w = \{\lambda w, \lambda \geq 0\}$. Let $\{q_n\}_{n \in \mathbb{N}} \subseteq r_w$ be a sequence as in Lemma 4.8. Then,

$$\partial_2 x(0, p_0, T)w = \lim_{n \rightarrow \infty} \frac{x(0, q_n, T) - x(0, p_0, T)}{|q_n - p_0|} = \lim_{n \rightarrow \infty} \frac{q_n - p_0}{|q_n - p_0|} = w.$$

The linearity of $\partial_2 x(0, p_0, T)$ and the arbitrariness of w show that $\partial_2 x(0, p_0, T)$ is, in fact, the identity. Take now any vector $v \in \mathbb{R}^2$ and consider the half-line r_v generated by v with origin in $(0, 0)$. Let $\{p_n\}_{n \in \mathbb{N}} \subseteq r_v$ be a sequence as in Lemma 4.8. Then we get

$$\begin{aligned} \partial_{22}x(0, p_0, T)[v, v] &= \lim_{n \rightarrow \infty} \frac{x(0, p_n, T) - x(0, p_0, T) - \partial_2 x(0, p_0, T)(p_n - p_0)}{\frac{1}{2}|p_n - p_0|^2} \\ &= \lim_{n \rightarrow \infty} \frac{p_n - p_0 - (p_n - p_0)}{\frac{1}{2}|p_n - p_0|^2} = 0, \end{aligned}$$

as desired. \square

Proof of Theorem 4.9. We have, for $t \in [0, T]$,

$$\begin{aligned} \partial_{22}x(0, p_0, t)v &= \\ &= \lim_{h \rightarrow 0^+} \frac{1}{h^2} \int_0^t \left[g(x(0, p_0 + hv, s)) - 2g(x(0, p_0, s)) + g(x(0, p_0 - hv, s)) \right] ds. \end{aligned}$$

Proceeding as in (9) we get

$$\begin{aligned} \partial_{22}x(0, p_0, T)[v, v] &= \\ &= \int_0^T \left[g'(p_0) \partial_{22}x(0, p_0, t)[v, v] + g''(p_0)[\partial_2x(0, p_0, s)v, x(0, p_0, s)v] \right] ds, \end{aligned}$$

Setting $\alpha(t) = \partial_2x(0, p_0, t)v$ and $\beta = \partial_{22}x(0, p_0, t)[v, v]$, we have that β solves the following Cauchy problem:

$$\begin{cases} \dot{\beta}(t) = g'(p_0)\beta(t) + g''(p_0)[\alpha(t), \alpha(t)], \\ \beta(0) = 0. \end{cases}$$

As we know, $\alpha(t) = e^{tg'(p_0)}v$. Thus

$$\beta(t) = \int_0^T e^{(T-s)g'(p_0)} g''(p_0) \left[e^{sg'(p_0)}v, e^{sg'(p_0)}v \right] ds.$$

The assertion follows from Lemma 4.10. \square

Remark 4.11. *Combining theorems 2.3 and 4.9 as in the argument of Lemma 3.2 we see that if p_0 is an isolated zero of g with the property that*

$$(14) \quad \exists v \in \mathbb{R}^2 : \int_0^T e^{(T-s)g'(p_0)} g''(p_0) \left[e^{sg'(p_0)}v, e^{sg'(p_0)}v \right] ds \neq 0,$$

then p_0 must be ejecting.

Condition (14) is difficult to use as it is, but it gets a considerably simpler form when $g'(p_0)$ is singular. In fact, one could prove that if $g'(p_0)$ is *not* singular, then (14) is never satisfied, in the sense that the integral in (14) is zero (the proof is based on the fact that when $g'(p_0)$ is not singular, but p_0 is T -resonant, $e^{tg'(p_0)}$ is necessarily a t -rotation matrix, so that the integrand has a particular form). Thus, the only interesting cases happen when $g'(p_0)$ is singular.

Let M, p_0 and g be as in Theorem 4.9. Suppose that $\ker g'(p_0) \neq \{0\}$. If $v \in \ker g'(p_0)$ then $e^{tg'(p_0)}v = v$. Thus, for such a v , equation (13) reads

$$\int_0^T e^{(T-s)g'(p_0)} g''(p_0)[v, v] ds = 0.$$

and condition (14) becomes

$$(15) \quad \exists v \in \ker g'(p_0) : \int_0^T e^{(T-s)g'(p_0)} g''(p_0)[v, v] ds \neq 0.$$

The situation is even simpler if $\ker g'(p_0) = \mathbb{R}^2$. In fact, in this case $e^{tg'(p_0)}$ is the identity for any t . Thus condition (14) reduces to

$$(16) \quad \exists v \in \mathbb{R}^2 : g''(p_0)[v, v] \neq 0.$$

We summarize the last discussion in the following

Theorem 4.12. *Let $M = \mathbb{R}^2$ and let p_0 be an isolated zero of g with nonzero index. Then, p_0 is ejecting for the set of T -pairs if there exists $v \in \ker g'(p_0)$ such that*

$$\int_0^T e^{(T-s)g'(p_0)} g''(p_0)[v, v] ds \neq 0.$$

If, in particular, $\ker g'(p_0) = \mathbb{R}^2$ it is sufficient to find $v \in \mathbb{R}^2$ such that $g''(p_0)[v, v] \neq 0$.

Example 4.13. *Let $g(x, y) = (x^3, y + x^2)$ and $T = 2\pi$. The only zero of g is the origin $(0, 0)$. We have*

$$g'(0, 0) = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix},$$

so that $(0, 0)$ is T -resonant but $i(g, (0, 0)) = 1$. Let us check if it is ejecting. Since $\det g'(0, 0) = 0$, we try to use condition (15). A quick computation shows that

$$e^{-sg'(0,0)} = \begin{pmatrix} 1 & 0 \\ 0 & e^{-s} \end{pmatrix}, \quad g''(0,0) \left[\begin{pmatrix} v_1 \\ v_2 \end{pmatrix}, \begin{pmatrix} v_1 \\ v_2 \end{pmatrix} \right] = \begin{pmatrix} 0 \\ 2v_2^2 \end{pmatrix}$$

Taking $v = \begin{pmatrix} 0 \\ 1 \end{pmatrix} \in \ker g'(0, 0)$, we get

$$\int_0^T e^{-sg'(0,0)} g''(0,0)[v, v] ds = \begin{pmatrix} 0 \\ 2(1 - e^{2\pi}) \end{pmatrix}.$$

Thus, by Theorem 4.12, $(0,0)$ is an ejecting zero of g . Figure 7 below shows some points sampled numerically from the set of starting points of (1) when g is as above and $f(t, x, y) = \sin t + 1$ (λ is on the vertical axis).

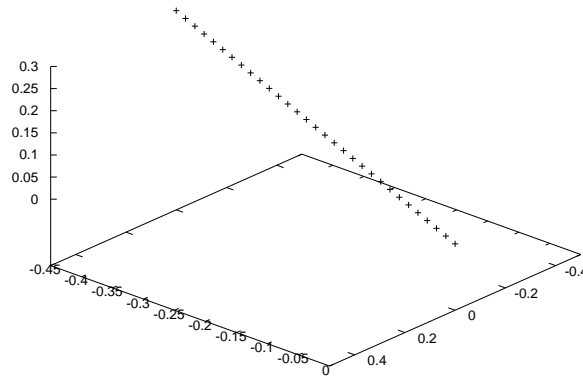


FIGURE 7. The set of starting points of Example 4.13.

Corollary 4.14. Let $g: M \rightarrow \mathbb{R}^k$ and $f: \mathbb{R} \times M \rightarrow \mathbb{R}^k$ be tangent to the closed boundaryless submanifold M of \mathbb{R}^k . Assume that g is C^2 , $g^{-1}(0)$ is compact and f is T -periodic in t . Assume also that

- (i) There are $n-1$ isolated zeros of g , p_1, \dots, p_{n-1} for which the conditions of Theorem 4.12 hold and such that

$$\sum_{i=1}^{n-1} i(g, p_i) \neq \deg(g, M);$$

- (ii) The unperturbed equation

$$\dot{x} = g(x)$$

does not admit unbounded connected sets of T -periodic solutions in $C_T(M)$.

Then, for $\lambda > 0$ sufficiently small, Equation (1) admits at least n geometrically distinct T -periodic solutions.

Remark 4.15. For the purpose of analyzing the notion of non- T -resonance we might have as well used an open set of \mathbb{R} and \mathbb{R}^2 . We have not done so, partly, for the sake of simplicity but, more importantly, because multiplicity results like the ones discussed in Section 3 require manifolds which are closed subsets of \mathbb{R}^k .

REFERENCES

- [1] L. Bisconti *Harmonic solutions to a class of differential-algebraic equations with separated variables*, Electronic Journal of Differential Equations, 2012, n. **2**, 1–15.
- [2] A. Boscaggin, A. Fonda, M. Garrione, *A multiplicity result for periodic solutions of second order differential equations with a singularity*, Nonlinear Anal. **75** (2012), no. 12, 4457–4470.
- [3] J. Cronin, *Differential Equations, Introduction and Qualitative Theory*, 2nd Ed., Pure and Appl. Math. 180, Marcel Dekker Inc., New York, 1994.
- [4] M. Furi and M. P. Pera, *Remarks on global branches of harmonic solutions to periodic ODE's on manifolds*, Boll. Un. Mat. Ital., **11-A** (1997), 709–722.
- [5] M. Furi, M. P. Pera and M. Spadini, *Forced oscillations on manifolds and multiplicity results for periodically perturbed autonomous systems*, Journal of computational and applied mathematics **113** (2000), 241–254.
- [6] M. Furi and M. Spadini, *On the set of harmonic solutions of periodic perturbations of differential equations on manifolds*, Journal of Nonlinear Analysis TMA, **9** (1997), 963–970.
- [7] M. Furi and M. Spadini, *Multiplicity of forced oscillations for the spherical pendulum*, Topological Methods in Nonlinear Analysis **11** (1998) 147–157.
- [8] T. Williams, C. Kelley et al., *Gnuplot*, <http://www.gnuplot.info>.
- [9] V. Guillemin and A. Pollack, *Differential Topology*, Prentice-Hall Inc., Englewood Cliffs, New Jersey, 1974.
- [10] M. W. Hirsch and S. Smale, *Differential equations, dynamical systems and linear algebra*. Academic Press, San Diego CA, 1974
- [11] P. Kunkel, V. Mehrmann, *Differential Algebraic Equations, analysis and numerical solutions*, EMS textbooks in Math., European Mathematical Society, Zürich 2006.
- [12] Lloyd N.G., *Degree Theory*, Cambridge Tracts in Math. 73, Cambridge University Press, Cambridge, 1978.
- [13] J. Mawhin, *Periodic oscillations of forced pendulum-like equations*, Ordinary and partial differential equations (Dundee, 1982), pp. 458–476, Lecture Notes in Math., 964, Springer, Berlin-New York, 1982.

- [14] J. Mawhin, *The implicit function theorem and its substitutes in Poincaré's qualitative theory of differential equations*, Stud. Hist. Philos. Sci. B Stud. Hist. Philos. Modern Phys. **47** (2014), 124–130.
- [15] J. W. Milnor, *Topology from the differentiable viewpoint*, Univ. press of Virginia, Charlottesville, 1965.
- [16] A. Pascoletti, M. Pireddu, F. Zanolin, *Multiple periodic solutions and complex dynamics for second order ODEs via linked twist maps*, The 8th Colloquium on the Qualitative Theory of Differential Equations, No. 14, 32 pp., Proc. Colloq. Qual. Theory Differ. Equ., 8, Electron. J. Qual. Theory Differ. Equ., Szeged, 2008.
- [17] C. Rebelo, F. Zanolin, *Multiple periodic solutions for a second order equation with one-sided superlinear growth*, Dynam. Contin. Discrete Impuls. Systems **2** (1996), no. 1, 1–27.
- [18] M. Spadini, *Harmonic solutions to perturbations of periodic separated variables ODEs on manifolds*, Electronic J. Diff. Equations **88** (2003) 1–11.
- [19] M. Spadini, *A note on topological methods for a class of Differential-Algebraic Equations*, Nonlinear Anal. **73** (2010), no. 4, 1065–1076.
- [20] A. J. Ureña, *Using closed orbits to bifurcate many periodic solutions for pendulum-type equations*, J. Math. Anal. Appl. **302** (2005), no. 2, 318–341.

(L. Bisconti and M. Spadini) DIPARTIMENTO DI MATEMATICA E INFORMATICA, UNIVERSITÀ DI FIRENZE,
VIA S. MARTA 3, 50139 FIRENZE, ITALY